Compression of Quantum Mixed State Sources

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A Classical Source of Information

- Discrete: produces sequences of letters.
- Letters belong to a finite alphabet \mathcal{X} .
- Memoryless: each letter is produced independently.

- Probability of letter a is P_x .
- Example: coin tossing with $\mathcal{X} = \{H, T\}$.
- ▶ Shannon Entropy: $-\sum_x P_x \log P_x$

A Quantum Source of Information

- Quantum letters are represented as unit-length vectors in \mathcal{H}_d .
- A qubit is a vector in \mathcal{H}_2 .
- **Example:** Alphabet $\mathcal{X} = \{0, 1, 2, 3\}$ mapped onto 4 qubits

$$\begin{split} |\psi_{0}\rangle &= \alpha_{0}|e_{0}\rangle + \beta_{0}|e_{1}\rangle \ |\psi_{1}\rangle = \alpha_{1}|e_{0}\rangle + \beta_{1}|e_{1}\rangle \\ |\psi_{2}\rangle &= \alpha_{2}|e_{0}\rangle + \beta_{2}|e_{1}\rangle \ |\psi_{3}\rangle = \alpha_{3}|e_{0}\rangle + \beta_{3}|e_{1}\rangle \end{split}$$

where $|e_0\rangle$ and $|e_1\rangle$ are the basis vectors of 2D space \mathcal{H}_2 :

$$|e_0
angle = \left[egin{array}{c} 0 \\ 1 \end{array}
ight] \qquad |e_1
angle = \left[egin{array}{c} 1 \\ 0 \end{array}
ight]$$

We will deal with (sequences of) qubits, WOLG.

The Density Matrix and Entropy

Source density matrix:

$$\rho = \sum_{\alpha \in \mathfrak{X}} \mathsf{P}_x \underbrace{|\psi_x\rangle \langle \psi_x|}_{\rho_x}.$$

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Von Neumann entropy of the source:

$$\begin{split} S(\rho) =& -\operatorname{Tr}\rho\log\rho\\ =& -\sum_i\lambda_i\log\lambda_i, \end{split}$$

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where λ_i are the eigenvalues of ρ .

THE MB EXAMPLE

$$\begin{split} \mathfrak{X} &= \{1, 2, 3\} \quad P_1 = P_2 = P_3 = 1/3 \\ \rho &= \frac{1}{3} |\psi_1\rangle \langle \psi_1| + \frac{1}{3} |\psi_2\rangle \langle \psi_2| + \frac{1}{3} |\psi_3\rangle \langle \psi_3| \\ &= \frac{1}{2} I \\ S(\rho) &= 1 \qquad \qquad |\psi_3\rangle = \begin{bmatrix} -1/2 \\ -\sqrt{3}/2 \end{bmatrix} \qquad \qquad |\psi_2\rangle = \begin{bmatrix} -1/2 \\ \sqrt{3}/2 \end{bmatrix} \end{split}$$

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Vector Sequences

Source vector-sequence (state)

$$|\Psi_{\mathbf{x}}\rangle = |\psi_{\mathbf{x}_1}\rangle \otimes |\psi_{\mathbf{x}_2}\rangle \otimes \cdots \otimes |\psi_{\mathbf{x}_n}\rangle, \qquad \mathbf{x}_i \in \mathfrak{X},$$

appears with probability $P_{\mathbf{x}} = P_{x_1} \cdot P_{x_2} \cdot \ldots \cdot P_{x_n}.$

- Typical states $|\Psi_{\mathbf{x}}\rangle \in \mathcal{H}^{2^n}$ correspond to typical sequences \mathbf{x} .
- ► There are approximately 2^{nH(P)} typical states.

Lossless Quantum Data Compression

- Source vector-sequence $|\Psi_{\mathbf{x}}\rangle$ is in \mathcal{H}^{2^n} , $(\mathbf{x}\in\mathfrak{X}^n)$
- Vector $|\Psi_{\mathbf{x}}\rangle$ is compressed and then reproduced as $|\widehat{\Psi_{\mathbf{x}}}\rangle$.
- Fidelity between $|\Psi_{\mathbf{x}}\rangle$ and $|\widehat{\Psi_{\mathbf{x}}}\rangle$:

 $\mathsf{F}(|\Psi_{\mathbf{x}}\rangle,|\widehat{\Psi_{\mathbf{x}}}\rangle) = |\langle \Psi_{\mathbf{x}}|\widehat{\Psi_{\mathbf{x}}}\rangle|^2$

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▶ For asymptotically lossless compression, the average fidelity

$$\overline{\mathbf{F}} = \sum_{\mathbf{X} \in \mathfrak{X}^n} \mathbf{P}(\mathbf{x}) \mathbf{F}(|\Psi_{\mathbf{x}}\rangle, |\widehat{\Psi_{\mathbf{x}}}\rangle)$$

should approach 1 as $n \to \infty$.

Typical States and Visible Compression

- Visible: the encoder Alice knows sequence x.
- She can compress with perfect fidelity the typical states.
- Instead of n qubits, she can transmit nH(P) bits.
- The decoder Bob prepares $|\Psi_{\mathbf{x}}\rangle$ as $|\widehat{\Psi_{\mathbf{x}}}\rangle$ for typical \mathbf{x} .

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- Can $|\Psi_{\mathbf{x}}\rangle$ be compressed to fewer than nH(P) qubits so that

- the compression is asymptotically lossless
- Alice does not know x
- Alice and Bob perform legal quantum operations

What Can be Done – Evolution (Reversible)

State ρ can be transformed to another state $\mathcal{E}(\rho)$ only by a physical process consistent with the lows of quantum theory:

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 $\mathcal{E}(\rho) = U \rho U^{\dagger}$ where $U U^{\dagger} = I$,

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unitary evolution:

$$\mathcal{E}(\rho) = U\rho U^{\dagger}$$
 where $UU^{\dagger} = I$,

completely positive, trace-preserving map:

$$\mathcal{E}(\rho) = \sum_k E_k \rho E_k^\dagger \quad \text{where} \quad \sum_k E_k^\dagger E_k = I.$$

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What Can be Done – Measurement (Irreversible)

Von Neumann:

- A set of pairwise orthogonal projection operators {Π_i}.
- They form a complete resolution of the identity: $\sum_{i} \Pi_{i} = I$.

• $|\psi_j\rangle$ is measured as $\Pi_i |\psi_j\rangle$ with probability $\langle \psi_j | \Pi_i | \psi_j \rangle$.

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- Positive Operator-Valued Measure (POVM):
 - ► Any set of positive-semidefinite operators {E_i}.
 - They form a complete resolution of the identity: $\sum_{i} E_{i} = I$.

• $|\psi_j\rangle$ is measured as $E_i |\psi_j\rangle$ with probability $\langle \psi_j | E_i | \psi_j \rangle$.

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► The No-Cloning Principle:

There is no physical process that leads to an evolution

 $|\varphi\rangle\otimes|s\rangle\rightarrow|\varphi\rangle\otimes|\varphi\rangle$

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where $|\phi\rangle$ is an arbitrary state and $|s\rangle$ is a fixed state.

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The No-Broadcasting Principle – generalization of no-cloning.

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- The No-Broadcasting Principle generalization of no-cloning.
- ► The No-Deleting Principle:

There is no physical process that leads to an evolution

$$|\phi\rangle\otimes|\phi\rangle\rightarrow|\phi\rangle\otimes|s\rangle$$

where $|\phi\rangle$ is an arbitrary state and $|s\rangle$ is a fixed state.

Typical Subspace

• Typical states $|\Psi_{\mathbf{x}}\rangle \in \mathcal{H}^{2^n}$ "live" in the typical subspace Λ_n .



The Typical Subspace Λ_n

We represent the source density matrix

$$\rho = \sum_{\alpha \in \mathfrak{X}} \mathsf{P}(x) |\psi_x\rangle \langle \psi_x|$$

in terms of its eigenvectors and eigenvalues:

$$\rho = \lambda_0 |\phi_0\rangle \langle \phi_0| + \lambda_1 |\phi_1\rangle \langle \phi_1|.$$

- Note that $\lambda = \{\lambda_0, \lambda_1\}$ is a PD on $\{0, 1\}$ and $\langle \phi_0 | \phi_1 \rangle = 0$.
- T_{λ}^{n} denotes the set of λ -typical sequences.
- Λ_n is the subspace spanned by $|\Phi_z\rangle$, $z \in \mathsf{T}^n_\lambda$.

Compression by Measurement

- \blacktriangleright Measurement is defined by $\Pi + \Pi^\perp = I_{2^n}$ where
 - $\Pi = \sum_{z \in T_{\lambda}^n} |\Phi_z\rangle \langle \Phi_z|$ is the projector to Λ_n .
 - $\Pi^{\perp} = \sum_{z \in \{0,1\}^n \setminus T_{\lambda}^n} |\Phi_z\rangle \langle \Phi_z|$ is the projector to Λ_n^{\perp} .

- State after measurement:
 - $\Pi \cdot |\Psi_x\rangle$ with probability $|\langle \Psi_x | \Pi | \Psi_x \rangle|^2$
 - $\Pi^{\perp} \cdot |\Psi_{\mathbf{x}}\rangle$ with probability $|\langle \Psi_{\mathbf{x}} | \Pi^{\perp} | \Psi_{\mathbf{x}} \rangle|^2$

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- State after measurement:
 - $\Pi \cdot |\Psi_x\rangle$ with probability $|\langle \Psi_x | \Pi | \Psi_x \rangle|^2$
 - $\Pi^{\perp} \cdot |\Psi_{\mathbf{x}}\rangle$ with probability $|\langle \Psi_{\mathbf{x}} | \Pi^{\perp} | \Psi_{\mathbf{x}} \rangle|^2$
- Expected probability of outcome $\Pi \cdot |\Psi_{\mathbf{x}}\rangle$:

$$\begin{split} \sum_{\mathbf{x}\in\mathfrak{X}^{n}} \mathsf{P}(\mathbf{x}) |\langle \Psi_{\mathbf{x}} | \Pi | \Psi_{\mathbf{x}} \rangle|^{2} \geqslant &-1 + 2\operatorname{Tr}(\Pi \rho^{\otimes n}) \\ = &-1 + 2\operatorname{Tr}\Big\{ \Big[\sum_{\mathbf{z}\in\mathsf{T}^{n}_{\lambda}} |\Phi_{z}\rangle \langle \Phi_{z}| \Big] \cdot \Big[\sum_{\mathbf{z}\in\{0,1\}^{n}} \lambda(z) |\Phi_{z}\rangle \langle \Phi_{z}| \Big] \Big\} \\ = &1 - 2\varepsilon_{n} \end{split}$$

- ► To a source letter x ∈ X corresponds quantum state |ψ_y⟩, y ∈ 𝔅, with probability W(y|x).
- Note that outputs are distributed as

$$Q(\mathbf{y}) = \sum_{\mathbf{x} \in \mathcal{X}} P(\mathbf{x}) W(\mathbf{y}|\mathbf{x}).$$

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The density matrix corresponding to x is

$$\rho_{x} = \sum_{\mathfrak{b} \in \mathfrak{Y}} W(\mathfrak{y}|x) |\psi_{\mathfrak{y}}\rangle \langle \psi_{\mathfrak{y}}|, \ x \in \mathfrak{X}.$$

Compression is asymptotically lossless when

$$\sum_{\mathbf{X}\in\mathfrak{X}^n}\mathsf{P}(\mathbf{x})\mathsf{F}(\mathbf{
ho}_{\mathbf{x}},\,\hat{\mathbf{
ho}}_{\mathbf{x}})
ightarrow 1$$
 as $\mathfrak{n}
ightarrow\infty$

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Produce sequences of sources, e.g., coins:



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A quantum example:

$$\begin{split} \rho_{1} &= \frac{2}{3} |\psi_{1}\rangle \langle \psi_{1}| + \frac{1}{3} |\psi_{2}\rangle \langle \psi_{2}| \\ \rho_{2} &= \frac{1}{3} |\psi_{2}\rangle \langle \psi_{2}| + \frac{2}{3} |\psi_{3}\rangle \langle \psi_{3}| \\ \rho &= \frac{1}{2} \rho_{1} + \frac{1}{2} \rho_{2} \\ &= \frac{1}{3} |\psi_{1}\rangle \langle \psi_{1}| + \frac{1}{3} |\psi_{2}\rangle \langle \psi_{2}| + \frac{1}{3} |\psi_{3}\rangle \langle \psi_{3}| \\ &= \frac{1}{2} I \\ & |\psi_{3}\rangle = \begin{bmatrix} -1/2 \\ -\sqrt{3}/2 \end{bmatrix}$$

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Distances Between Density Matrices ρ and σ

Uhlman fidelity:

$$F(\sigma, \omega) = \left\{ \mathsf{Tr} \big[(\sqrt{\sigma} \omega \sqrt{\sigma})^{1/2} \big] \right\}^2.$$

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► Trace distance:

$$D(\sigma, \omega) = \frac{1}{2} \operatorname{Tr} |\sigma - \omega|,$$

|A| denotes the positive square root of $A^{\dagger}A$.

► $1 - F(\sigma, \omega) \leq D(\sigma, \omega) \leq \sqrt{1 - F(\sigma, \omega)^2}$

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- ► $1 F(\sigma, \omega) \leq D(\sigma, \omega) \leq \sqrt{1 F(\sigma, \omega)^2}$
- Frobenius (Hilbert-Schmidt)?

Distances Between PD's - An Example

$$\blacktriangleright \mathcal{A}_{N} = \{a_{1}, \ldots, a_{N}\}$$

▶
$$N = 2^{K}$$
 and $n = 2^{k}$, with $k/K = c < 1$.

Distributions P and Q:

$$P(a_i) = \begin{cases} 1/n, & 1 \leqslant i \leqslant n \\ 0 & n+1 \leqslant i \leqslant N \end{cases} \text{ and } Q(a_i) = \frac{1}{N}$$

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►
$$Q(\{a_{n+1}, \dots, a_N\}) \rightarrow 1 \text{ as } k, K \rightarrow \infty.$$

▶ $\frac{1}{2}\sum_i |P(a_i) - Q(a_i)| \rightarrow 1 \text{ and } \sum_i |P(a_i) - Q(a_i)|^2 \rightarrow 0.$

Sources of Sources

Produce sequences of sources, e.g., coins:



• Example:
$$P(C_1) = P(C_2) = 1/2$$
 and $w = 2/3$.

Sequences:

C1	C_1	C ₂	C_1	C ₂	C ₂
н	Н	Н	Т	Т	Т
Н	Т	Н	Н	Т	Т
Н	т	т	Н	т	н

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Sequences of Sources

▶ For each source (coin) x in X,

we have a probability distribution $W(\cdot|\mathbf{x})$ over \mathcal{Y} of letters (faces).

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- ► For each sequence of coins x in Xⁿ, we have a probability distribution Wⁿ(·|x) over Yⁿ.
- Alice has coins x and sends N bits to Bob.
- Bob prepares faces \mathbf{y} with probability $\widehat{W}^n(\mathbf{y}|\mathbf{x})$.
- $W^n(\cdot|\mathbf{x})$ is reproduced as $\widehat{W}^n(\cdot|\mathbf{x})$.

Compression by Sending Classical Information

 \blacktriangleright Fidelity between $W^n(\cdot|\mathbf{x})$ and $\widehat{W}^n(\cdot|\mathbf{x})$

$$F(W^{n}(\cdot|\mathbf{x}),\widehat{W}^{n}(\cdot|\mathbf{x})) = \sum_{\mathbf{y}\in\mathcal{Y}^{n}} \sqrt{W^{n}(\mathbf{y}|\mathbf{x}) \cdot \widehat{W}^{n}(\mathbf{y}|\mathbf{x})}$$

is known as the Bhattacharyya-Wooters overlap.

Compression fidelity

$$\sum_{\mathbf{x}\in\mathfrak{X}^{n}}\mathsf{P}(\mathbf{x})\mathsf{F}\big(\mathsf{W}^{n}(\cdot|\mathbf{x}),\widehat{W}^{n}(\cdot|\mathbf{x})\big)$$

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should approach 1 as $n \to \infty$.

Compression Algorithm for a Typical x

► Alice and Bob

- have identical random number generators
- which they use to form a list of N₁ typical ys.
- If N_l > 2^{nI(P,W)}, then with high probability there will be at least one y on the list which is conditionally typical with respect to x.

- Alice sends $\log N_1$ bits to Bob identifying y.
- Compression rate $\log N_l/n$ approaches I(P, W).

Proof Idea

- ► Wⁿ(·|x) is roughly a uniform distribution over ys that are conditionally typical given x.
- For a typical x, there are about $2^{nH(W/P)}$ such ys.
- ► These ys are typical.
- There are about $2^{nH(Q)}$ typical **y**s.
- A randomly chosen y will be conditionally typical with respect to any typical x with probability of about

$$\frac{2^{nH(W/P)}}{2^{nH(Q)}} = \frac{1}{2^{nI(P,W)}}.$$

A Related Problem

For each Alice's sequence $C_{\mathbf{x}}$ of coins,

Bob prepares a predetermined sequence y(x) of faces such that

$$\bar{\mathsf{F}} = \sum_{\mathbf{x} \in \mathfrak{X}^{n}} \mathsf{P}(\mathbf{x}) \mathsf{F}_{\mathfrak{X} \times \mathfrak{Y}} \big(\mathsf{P}_{\mathbf{x}} \mathsf{W}(\cdot | \cdot), \mathsf{P}_{\mathbf{x}, \mathfrak{Y}(\mathbf{x})} \big)$$

$$\begin{split} \mathsf{F}_{\mathfrak{X}\times\mathfrak{Y}}\big(\mathsf{P}_{\mathbf{x}}W(\cdot|\cdot),\mathsf{P}_{\mathbf{x},\mathbf{y}\,(\mathbf{x})}\big) &= \Bigl[\sum_{(\mathbf{x},\mathbf{y})\in\mathfrak{X}\times\mathfrak{Y}}\sqrt{\mathsf{P}_{\mathbf{x}}(\mathbf{x})W(\mathbf{y}|\mathbf{x})\cdot\mathsf{P}_{\mathbf{x},\mathbf{y}\,(\mathbf{x})}(\mathbf{x},\mathbf{y})}\Bigr]^{2} \\ &= \Bigl[\sum_{(\mathbf{x},\mathbf{y})\in\mathfrak{X}\times\mathfrak{Y}}\frac{1}{n}\sqrt{\mathsf{N}(\mathbf{x}|\mathbf{x})W(\mathbf{y}|\mathbf{x})\cdot\mathsf{N}(\mathbf{x},\mathbf{y}|\mathbf{x},\mathbf{y}(\mathbf{x}))}\Bigr]^{2}. \end{split}$$

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How large is Bob's codebook?

The Original Quantum Problem

• The compression rate can not go below the Holevo quantity χ :

$$\chi = \underbrace{-\operatorname{Tr} \rho \log \rho}_{S(\rho)} - \sum_{\alpha \in \mathcal{X}} P(\alpha) S(\rho_{\alpha})$$

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- I(P, W) is achievable by sending classical information.
 The proof uses
 - 1. the equivalence of the Uhlman fidelity and the trace distance
 - 2. the strong convexity of the trace distance:

$$D\left(\sum_{i} p_{i}\omega_{i}, \sum_{i} q_{i}\sigma_{i}\right) \leqslant D(\{p_{i}\}, \{q_{i}\}) + \sum_{i} p_{i}D(\omega_{i}, \sigma_{i}).$$

3. the method of types

The Original Quantum Problem

• The compression rate can not go below the Holevo quantity χ :

$$\chi = \underbrace{-\operatorname{Tr} \rho \log \rho}_{S(\rho)} - \sum_{\alpha \in \mathcal{X}} P(\alpha) S(\rho_{\alpha})$$

- I(P, W) is achievable by sending classical information.
 The proof uses
 - 1. the equivalence of the Uhlman fidelity and the trace distance
 - 2. the strong convexity of the trace distance:

$$D\left(\sum_{i} p_{i}\omega_{i}, \sum_{i} q_{i}\sigma_{i}\right) \leq D(\{p_{i}\}, \{q_{i}\}) + \sum_{i} p_{i}D(\omega_{i}, \sigma_{i}).$$

- 3. the method of types
- Can the gap be closed by qubits or bits?